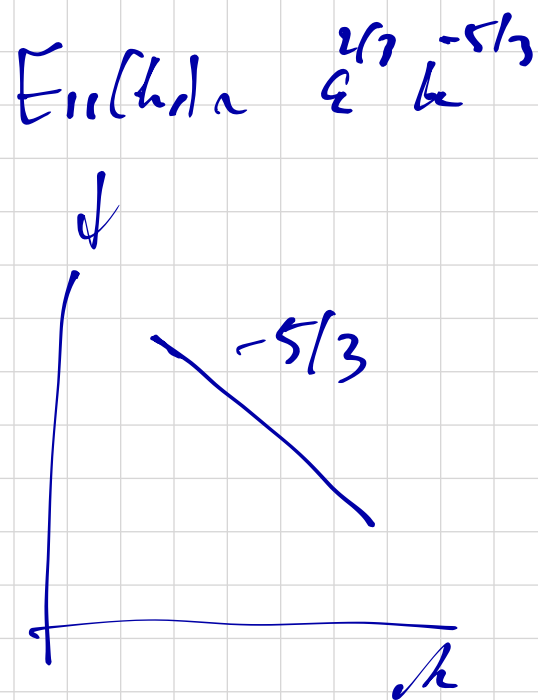
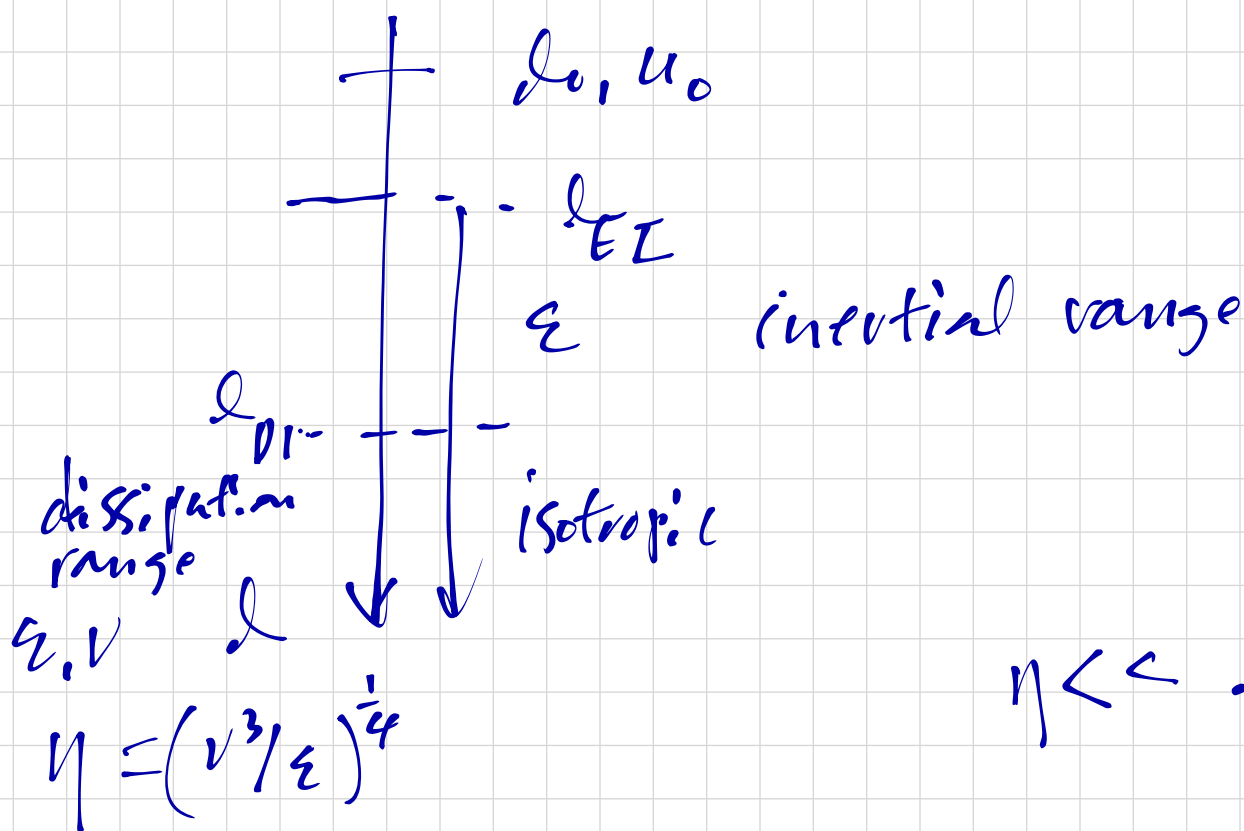


10/16/2018

Recap Energy Cascade



Kolmogorov (1941b)

$$P_{ij}(r) = \underset{\uparrow}{A} \delta_{ij} + \underset{\uparrow}{B} r_i r_j$$

P_{LL}, P_{NN}

In particular, inertial range

$$P_{ij} = C_2 (\epsilon/r)^{2/3} \left(\frac{4}{3} \delta_{ij} - \frac{1}{3} \frac{r_i r_j}{r^2} \right) \leftarrow$$

Aches Exp.

$$Re_\delta = 3.6 \times 10^6$$

$$\delta = 1.09 \text{ mm}$$

$$\eta = 0.09 \text{ mm}$$

$$\delta/\eta \sim 12,000$$

$$\frac{P_{11}}{(\epsilon r)^{2/3}} = C_2$$

$$\frac{P_{22}}{(\epsilon r)^{2/3}} = \frac{P_{33}}{(\epsilon r)^{2/3}} = \frac{4}{3} C_2$$

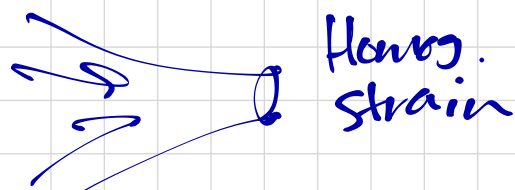
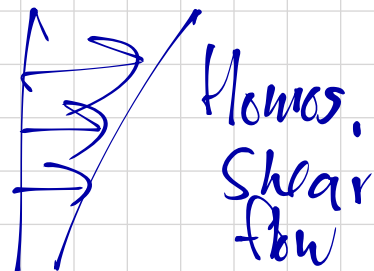
Do K-4 here
(see notes)

Evolution of the Energy Spectrum

(Spectral view of energy cascade)

See Hinze (1975), Monin & Yaglom (1975)

For homogeneous turbulence



$$\frac{\partial}{\partial t} E(k, t) = \underbrace{P_k(k, t)}_{\text{Production}} - \underbrace{\frac{\partial}{\partial k} T_k(k, t)}_{\text{Transfer}} - \underbrace{2\nu k^2 E(k, t)}_{\substack{D(k, t) \\ \text{Dissipation}}}$$

$$P_k(k, t) = \int_{k_a}^{k_b} P_k(k, t) dk \quad \text{contribution to the production from } (k_a, k_b)$$

$$P = \int_0^{\infty} P_k(k, t) dk \quad \text{total production}$$

T_k : spectral energy transfer rate, which is net rate at which energy is transferred from the lower k to larger k

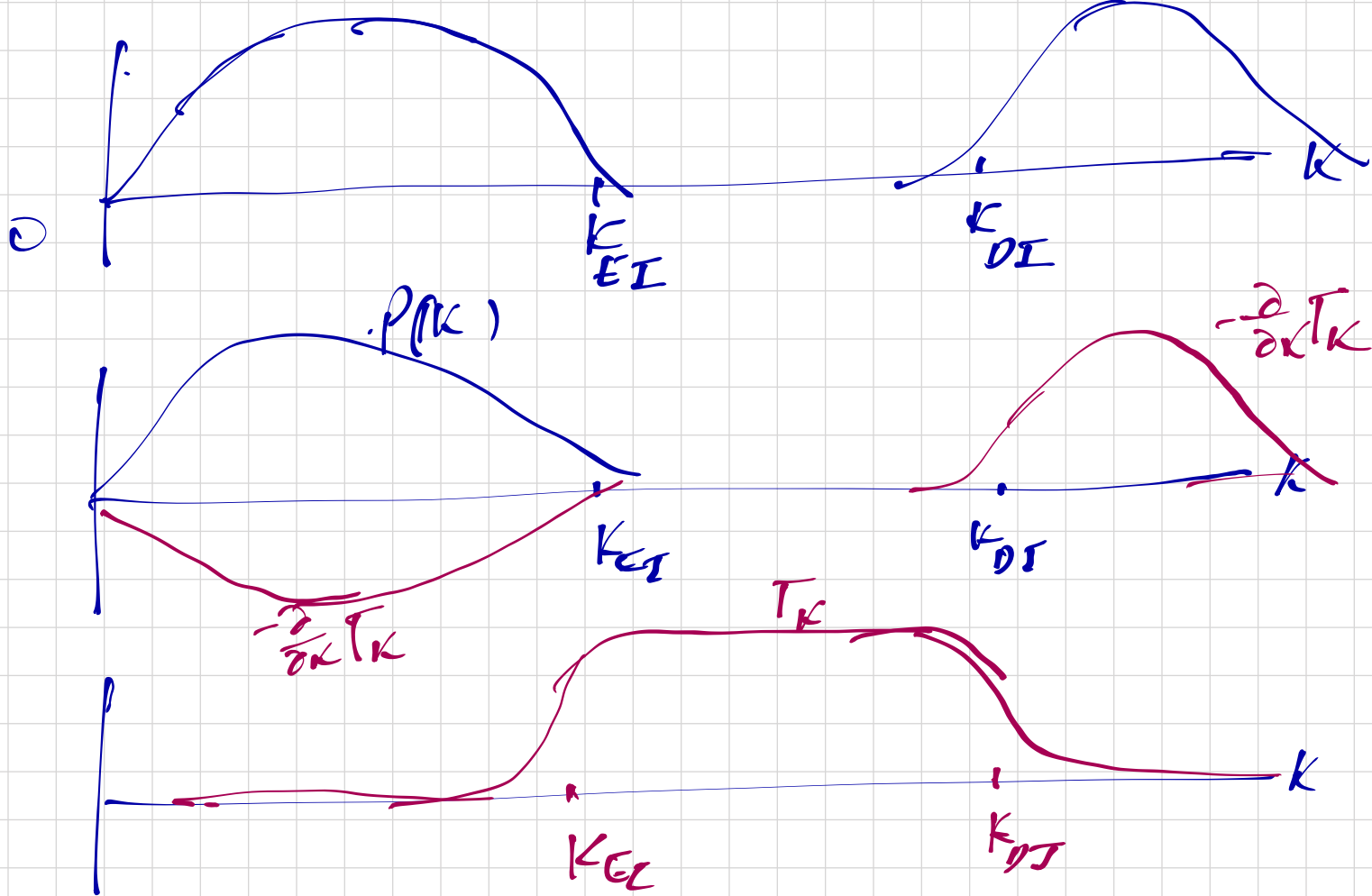
$$\int_{k_a}^{k_b} - \frac{\partial}{\partial k} T_k(k, t) dk = T_k(k_a) - T_k(k_b)$$

$$D: \text{Dissipation} \quad D(k, t) = 2\nu k^2 E(k, t)$$

$$\int_0^{\infty} D(k, t) dk = \epsilon \quad \text{total dissipation}$$

$E(k)$

$D(k)$



Energy Containing range $(0, k_{EI})$

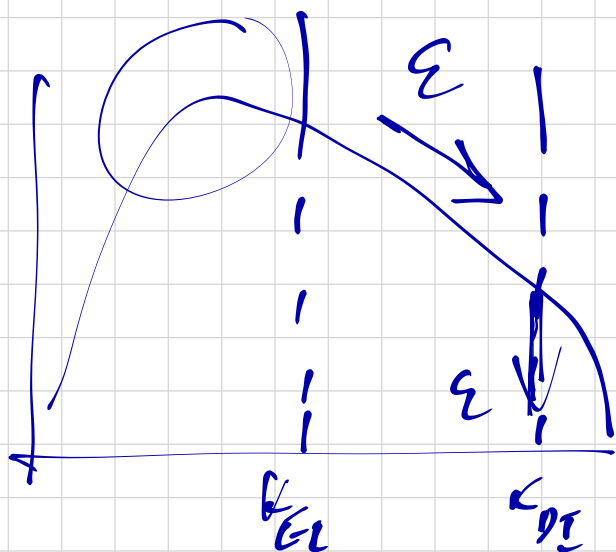
$$\int_0^{k_{EI}} E(k) dk \sim \int_0^\infty E(k) dk = k \quad \checkmark TKE$$

$$(0, k_{EI}) \quad \frac{dk}{dt} \sim P - T_{EI}$$

$$(k_{EI}, k_{DI}) \quad 0 \sim T_{EI} - T_{DI}$$

$$(k_{DI}, \infty) \quad 0 \sim T_{DI} - \varepsilon$$

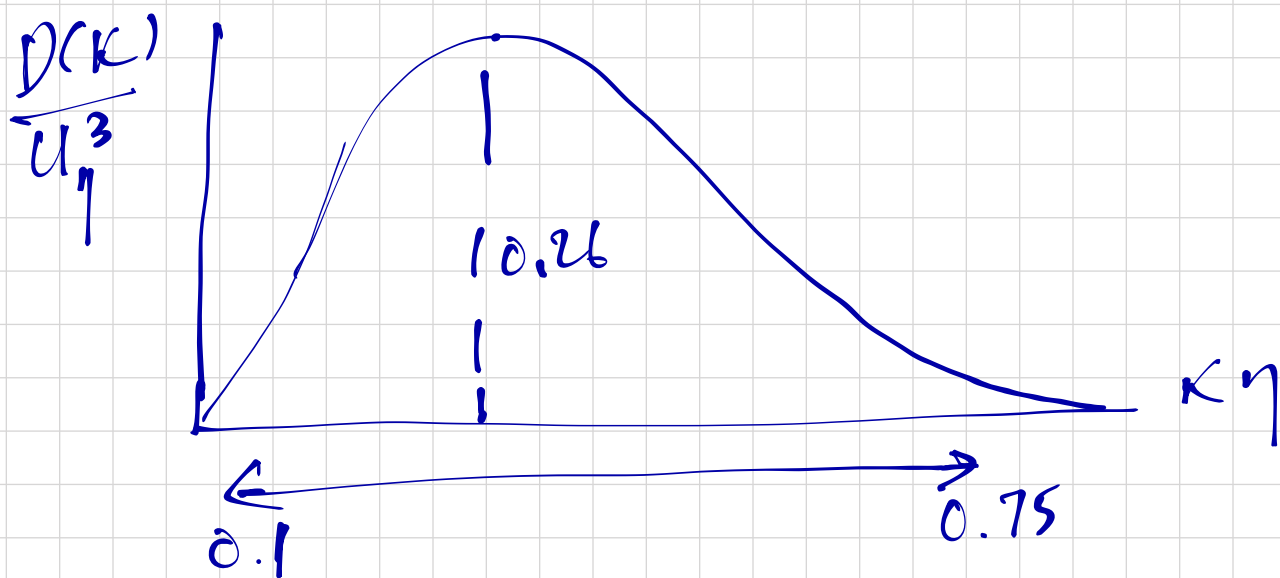
$$\sum \frac{dk}{dt} = P - \varepsilon$$



Dissipation Spectrum

$$\epsilon = \int_0^{\infty} p(k) dk$$

$\uparrow$
 $2\nu k^2 E(k)$



$$k\eta \approx 0.1$$

$$\frac{2\pi}{\lambda} \eta \approx 0.1$$

$$\frac{2\pi}{\lambda} \eta \approx 0.75$$

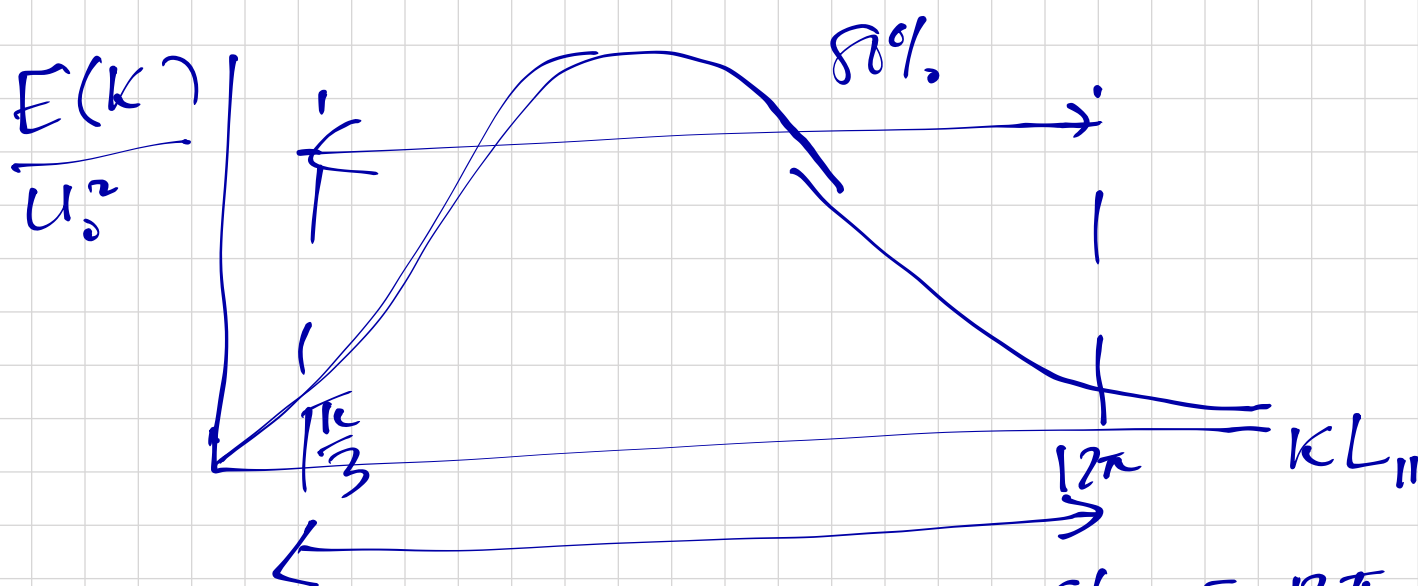
$$\frac{\lambda}{\eta} \approx 60$$

$$\frac{\lambda}{\eta} = 8$$

$$\lambda_{DI} \approx 60\eta$$

Energy Containing Range

$$k = \int_0^{\infty} E(k) dk$$



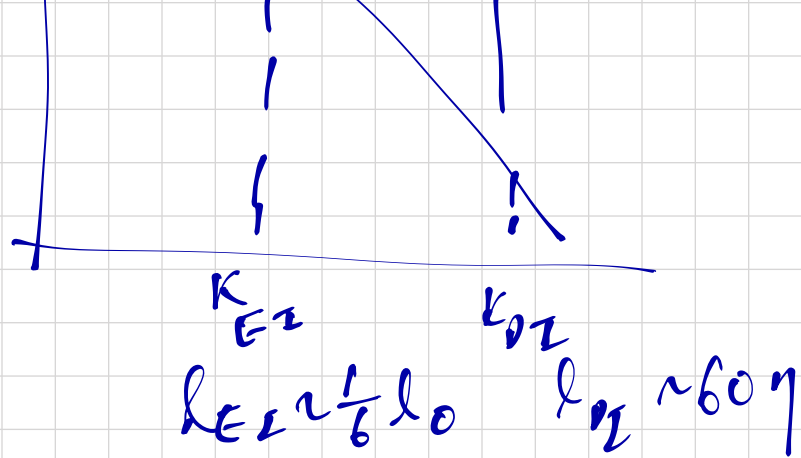
$$kL_\eta = 12\pi$$

$$\frac{2\pi}{\lambda} L_\eta = 12\pi \rightarrow \frac{\lambda}{L_\eta} = \frac{1}{6}$$

$$kL_\eta = \frac{\pi}{3}$$

$$\rightarrow \lambda = 6L_\eta$$





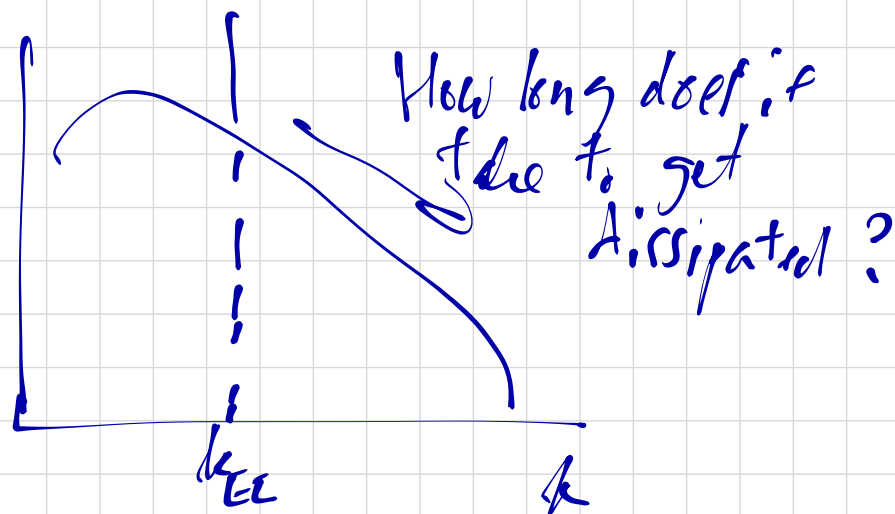
Note

1. Velocity scale for energy containing eddies $\sqrt{k}$

2. Time scale $L_{\eta} / \sqrt{k} \rightarrow L_{\eta} = k^{3/2} / \epsilon$
 Time scale $\tau = k / \epsilon \rightarrow$ Eddy turn over time

3. Cascade time scale?

$\sim \frac{1}{10} \tau$
 (See notes)



Limitations, Shortcomings and Refinements
 (Kolmogorov 1962)

1. How large should R_{η} ?
2. Energy cascade; reverse cascade?
 locally, instantaneously
3. Non-local interaction (i.e. between large and small)
 scale
4. $D_{\alpha j}$ Higher-order statistics tend to be deviated from Kolmogorov hypothesis

5. Intermittency

6. Refined Similarity ϵ

ϵ_r : local value of ϵ , average over $r \ll L$

$\rightarrow$ Stats $\Delta_r u$ depend on ϵ_r, ν

$\rightarrow$ JFM, 1995 JFM 297 Stolovitzky et al.

$$L, \lambda, \eta$$

$$L > \lambda > \eta$$

$$\epsilon \sim \left(\frac{\partial u_i}{\partial x_j} \right)^2 \sim \left(\frac{u_j}{\eta} \right)^2 \sim \left(\frac{u}{\lambda} \right)^2$$

$$\omega_i \sim \frac{\partial u}{\partial y} \sim \frac{u}{\lambda}$$

What about $\frac{\partial \omega_i}{\partial x_j}$?

Vorticity Equation

$$\frac{D}{Dt} \frac{\omega_i \omega_i}{2} = \dots$$

$$\omega_i \omega_j S_{ij}$$

vortex stretching

$$-\nu \frac{\partial \omega_i}{\partial x_j} \frac{\partial \omega_i}{\partial x_j}$$

viscous dissipation

For high Re

$$\omega_i \omega_j S_{ij} \sim \nu \frac{\partial \omega_i}{\partial x_j} \frac{\partial \omega_i}{\partial x_j}$$

$$S_{ij} = O\left(\frac{u}{\lambda}\right), \quad \omega = O\left(\frac{u}{\lambda}\right)$$

$$\frac{\partial \omega_i}{\partial x_j} ?$$

$$\left(\frac{u}{\lambda}\right)^3 = \nu \left(\frac{u}{\lambda}\right)^2 \frac{1}{l^2}$$

$$l = \left(\frac{\nu \lambda}{u}\right)^{\frac{1}{2}}$$

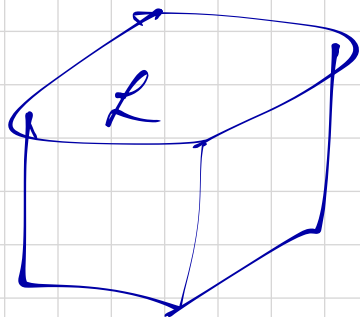
$$= \left(\frac{\nu^2 \lambda^2}{u^2}\right)^{\frac{1}{4}} = \left[\frac{\nu^3}{\nu \left(\frac{u}{\lambda}\right)^2}\right]^{\frac{1}{4}}$$

$$= \left(\frac{1^3}{2} \right)^{\frac{1}{4}} = \eta$$

Use the Kolmogorov scale to estimate the derivative of ω !

Fourier Modes

$$\underline{u}(\underline{x}, t) = \sum_{\underline{k}} \hat{\underline{u}}(\underline{k}, t) e^{i \underline{k} \cdot \underline{x}}$$



$L \gg L_{11}$

$$\text{Let } g(\underline{x}) = \sum_{\underline{k}} e^{i \underline{k} \cdot \underline{x}} \hat{g}(\underline{k})$$

$$\mathcal{F}\{g(\underline{x})\} = \hat{g}(\underline{k})$$

$$\mathcal{F}\left\{\frac{\partial}{\partial x_j} g(\underline{x})\right\} = i k_j \hat{g}(\underline{k})$$

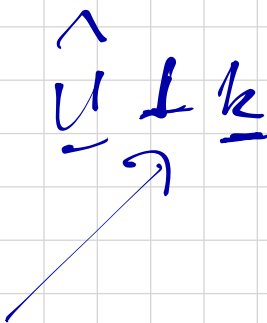
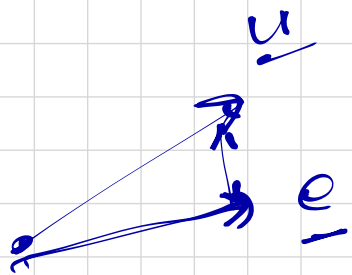
$$\mathcal{F}\{\nabla^2 g\} = -k^2 \hat{g}$$

Conti

$$\frac{\partial u_i}{\partial x_i} = 0$$

$$\mathcal{F}\left\{\frac{\partial u_i}{\partial x_i}\right\} \rightarrow i k_j \hat{u}_j = 0$$

$$\underline{k} \cdot \hat{\underline{u}} = 0$$



$$\underline{u}'' = (\underline{u} \cdot \underline{e}) \underline{e}$$

$$\underline{u}^\perp = \underline{u} - \underline{u}''$$

$$u_i'' = (u_j e_j) e_i$$

$$u_i^\perp = u_i - (u_j e_j) e_i$$

$$= u_j \delta_{ij} - u_j e_j e_i$$

$$= u_j (\delta_{ij} - e_i e_j)$$

$$\underline{e} = \underline{k} / |\underline{k}|$$

$$d_{ij} = \frac{k_i k_j}{k^2}$$

$$P_{ij}$$

$$P_{ij} = \left(\delta_{ij} - k_i k_j / k^2 \right)$$

Projection Tensor

$y \{ N-S \}$

$$y \left\{ \frac{\partial u_i}{\partial t} + \underbrace{\frac{\partial}{\partial x_j} u_i u_j}_{G_i} = - \frac{\partial p}{\partial x_i} + \nu \frac{\partial^2 u_i}{\partial x_j \partial x_j} \right\}$$

$$\frac{\partial}{\partial t} \hat{u}_i + \hat{G}_i = -i k_i \hat{p} - \nu k^2 \hat{u}_i$$

Poisson equ. for pressure

$$k^2 \hat{p} = i k_j \hat{G}_j$$

$$\hat{p} = i \frac{k_j}{k^2} \hat{G}_j$$

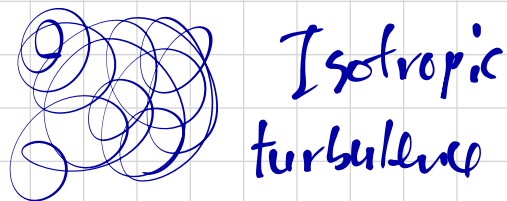
$$-i k_i \hat{p} = \frac{k_i k_j}{k^2} \hat{G}_j = \hat{G}_i^{\parallel}$$

$$\frac{\partial}{\partial t} \hat{u}_i + \hat{G}_i = \hat{G}_i^{\parallel} - \nu k^2 \hat{u}_i$$

$$\frac{\partial \hat{u}_i}{\partial t} = - \underbrace{\hat{G}_i^{\perp} + \hat{G}_i^{\parallel}}_{\hat{G}_i} - \nu k^2 \hat{u}_i$$

$$\left(\frac{\partial}{\partial t} + \nu k^2 \right) \hat{u}_i = - \hat{G}_i^{\perp} = - \underbrace{\left(\delta_{ij} - \frac{k_i k_j}{k^2} \right)}_{P_{ij}} \hat{G}_j$$

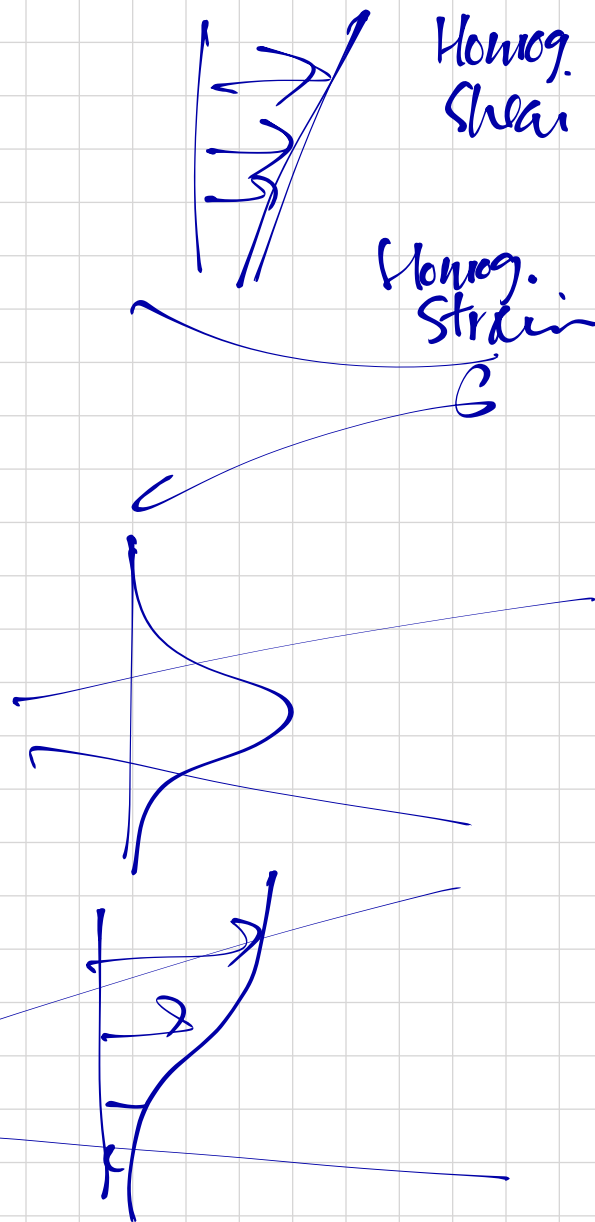
$$\frac{D}{Dt} k = P + T - \epsilon$$



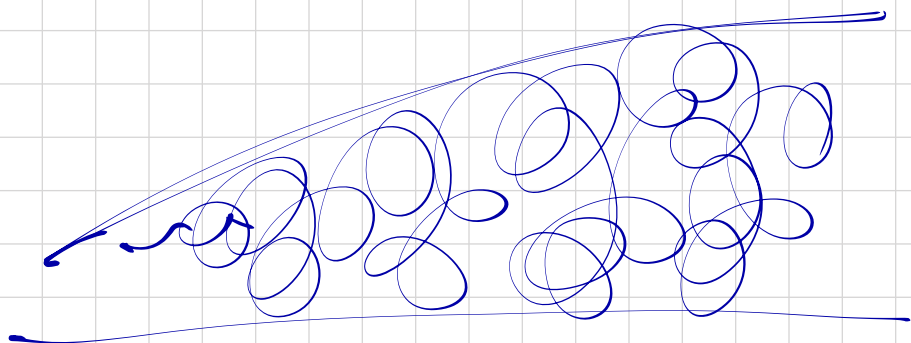
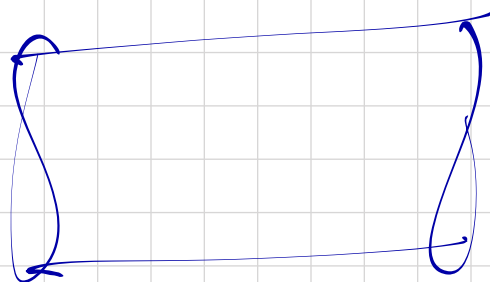
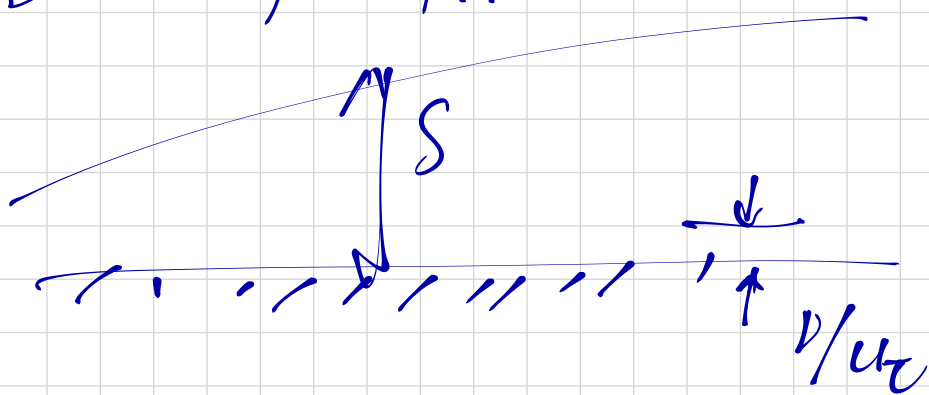
$$\frac{D}{Dt} k = P + T - \epsilon$$

Inhomog. Flows

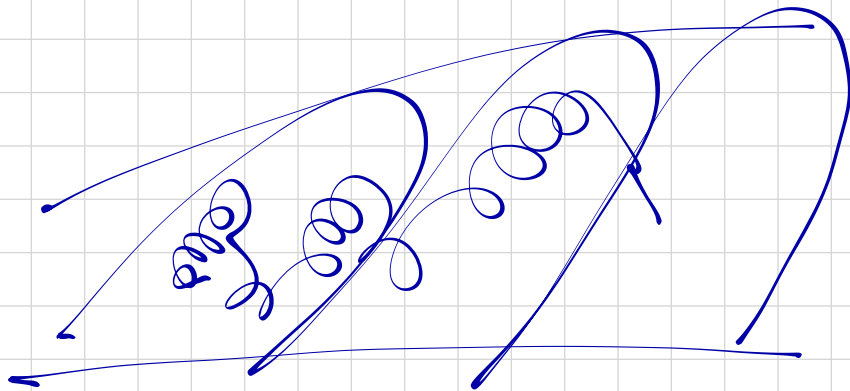
Free-Shear Flows



Boundary Layer Flow



Old View



New View